#### Multi-Scale Modeling of Spin Dynamics in Molecular Semi-Conductors

#### **Erik R. McNellis**

Research Team Leader Johannes Gutenberg Universität Mainz



#### April 10th, 2019 ITN SEPOMO | NWE4 Mons, Belgium







European Research Council



# Who We Are - ERC Synergy Grant



- ‣ ERC Synergy Grant focused on **organic spintronics**
- Interdisciplinary, joining theory / experiment / physics / chemistry / materials science



**European Research Council** 

#### ‣ PIs

- ‣ H. Sirringhaus, Cambridge
- ‣ J. Sinova, JGU Mainz
- ‣ I. McCulloch, Imperial College
- ‣ J. Wunderlich, Hitachi Cambridge
- ‣ Outside Synergy Grant, groups of
	- ‣ D. Andrienko, MPIP Mainz
	- ‣ D. Beljonne, University of Mons











#### Who We Are - Theory Team





Prof. Sergei A. Egorov



Dr. Reza Mahani



M. Sc. Uday Chopra



M. Sc. Sebastian Müller



https://www.sinova-group.physik.uni-mainz.de/research/organic-spintronics/

#### Why First-Principles, Multi-Scale Modeling?





- 1. Z. H. Xiong, D. Wu, Z. V. Vardeny and J. Shi, Nature 427, 821 (2004)
- 2. S. W. Jiang, S. Liu, P. Wang, Z. Z. Luan et al, Phys. Rev. Lett. 115, 086601 (2015)



### Why First-Principles, Multi-Scale Modeling?



- ‣ 'Fruit-fly' example: characteristic Alq3 spin dynamics **varies hugely** depending on e.g.
	- ‣ morphology
	- temperature
	- ‣ spin (charge) density





Molecular vibrations<sup>3</sup>

- 1. Z. H. Xiong, D. Wu, Z. V. Vardeny and J. Shi, Nature 427, 821 (2004)
- 2. S. W. Jiang, S. Liu, P. Wang, Z. Z. Luan et al, Phys. Rev. Lett. 115, 086601 (2015)
	- 3. L. Nuccio, M. Willis, L. Schulz, S. Fratini et al, Phys. Rev. Lett. 110, 216602 (2013)



### Why First-Principles, Multi-Scale Modeling?



- ‣ 'Fruit-fly' example: characteristic Alq3 spin dynamics **varies hugely** depending on e.g.
	- ‣ morphology
	- ‣ temperature
	- ‣ spin (charge) density



Spin exchange2

- ‣ Need modeling **consistently accurate** across spintronic device designs / operating regimes
- ‣ Phenomenological models struggle



Molecular vibrations<sup>3</sup>

- 
- 1. Z. H. Xiong, D. Wu, Z. V. Vardeny and J. Shi, Nature 427, 821 (2004)
- 2. S. W. Jiang, S. Liu, P. Wang, Z. Z. Luan et al, Phys. Rev. Lett. 115, 086601 (2015)
- 3. L. Nuccio, M. Willis, L. Schulz, S. Fratini et al, Phys. Rev. Lett. 110, 216602 (2013)



- ‣ Molecular / organic semi-conductors characterized by
	- lower order / crystallinity, charge mobility ( / hopping frequency?)
	- weaker local fields, spin-orbit coupling (more 'orbits' relevant, 'SOCs'?)
	- ‣ larger morphological variation / anisotropy



- ‣ Molecular / organic semi-conductors characterized by
	- lower order / crystallinity, charge mobility (/ hopping frequency?)
	- ‣ weaker local fields, spin-orbit coupling (more 'orbits' relevant, 'SOCs'?)
	- larger morphological variation / anisotropy
- ‣ For hopping charge transport, no inter-system crossings **five spin relaxation mechanisms**
- 1. Spin dipole:
	- orients spins (anti-) parallel for (perpendicular) parallel separation vector **R**
	- ‣ ignored in solid state, **can** matter in organics (short **R**)
	- ‣ modeled classically







#### Molecular Spin Relaxation Mechanisms

- ‣ Molecular / organic semi-conductors characterized by
	- lower order / crystallinity, charge mobility (/ hopping frequency?)
	- weaker local fields, spin-orbit coupling (more 'orbits' relevant, 'SOCs'?)
	- larger morphological variation / anisotropy
- ‣ For hopping charge transport, no inter-system crossings **five spin relaxation mechanisms**
- 1. Spin dipole:
	- orients spins (anti-) parallel for (perpendicular) parallel separation vector **R**
	- ‣ ignored in solid state, **can** matter in organics (short **R**)
	- ‣ modeled classically
- 2. Spin exchange: spontaneous inversion of neighboring spins
- 3. Hyperfine fields: due to electronic / nuclear spin interaction
- 4. Hop-flipping: scattering between mixed spin states  $\rightarrow$  spin flip
- 5. Thermal (spin-phonon coupling): phonon scattering + SOC







#### Molecular Spin Relaxation Mechanisms

- ‣ Molecular / organic semi-conductors characterized by
	- lower order / crystallinity, charge mobility (/ hopping frequency?)
	- weaker local fields, spin-orbit coupling (more 'orbits' relevant, 'SOCs'?)
	- ‣ larger morphological variation / anisotropy
- ‣ For hopping charge transport, no inter-system crossings **five spin relaxation mechanisms**
- 1. Spin dipole:
	- orients spins (anti-) parallel for (perpendicular) parallel separation vector **R**
	- ‣ ignored in solid state, **can** matter in organics (short **R**)
	- ‣ modeled classically
- 2. Spin exchange: spontaneous inversion of neighboring spins
- 3. Hyperfine fields: due to electronic / nuclear spin interaction
- Hop-flipping: scattering between mixed spin states  $\rightarrow$  spin flip
- 5. Thermal (spin-phonon coupling): phonon scattering + SOC





Least explored in this context



*"The underlying physical laws necessary for the mathematical theory of a large part of physics and the whole of chemistry are thus completely known, and the difficulty is only that the exact application of these laws leads to* **equations much too complicated to be soluble***. It therefore becomes desirable that*  **approximate practical methods** *of applying quantum mechanics should be developed, which can lead to an explanation of the main features of complex atomic systems*  **without too much computation.***"* 



1. P. A. M. Dirac, Proc. R. Soc. Lond. A 123, 714 (1929)



- ‣ A quantum wave function contains a **lot of information**. Ideally, we
	- ✓ want **all necessary** information (accuracy)
	- ✓ do **not** want to **filter** for specific information (transferability)
	- ✓ do **not** want to calculate **too much** (scalability)



#### First-Principles Modeling in a Nutshell



- ‣ A quantum wave function contains a **lot of information**. Ideally, we
	- ✓ want **all necessary** information (accuracy)
	- ✓ do **not** want to **filter** for specific information (transferability)
	- ✓ do **not** want to calculate **too much** (scalability)
- ✗ *Contradiction*



**Transferability**





- ‣ A quantum wave function contains a **lot of information**. Ideally, we
	- ✓ want **all necessary** information (accuracy)
	- ✓ do **not** want to **filter** for specific information (transferability)
	- ✓ do **not** want to calculate **too much** (scalability)
- ✗ *Contradiction*



**Transferability**

- ‣ Density Functional Theory (DFT): correction potential for classical mean-field
- ‣ Strikes good **balance** between all three popular method







- ‣ **Perfect** theoretical material model:
	- ✓ From atomic- to material-relevant scale
	- ✓ No empiricism
	- ✓ Accurate, transferable
	- ✓ Computable



#### Multi-Scale Modeling of Materials

- ‣ **Perfect** theoretical material model:
	- ✓ From atomic- to material-relevant scale
	- ✓ No empiricism
	- ✓ Accurate, transferable
	- ✓ Computable
- ‣ **First-Principles Modeling:** 
	- ‣ Accurate for single component, impractical for material







#### Multi-Scale Modeling of Materials

- ‣ **Perfect** theoretical material model:
	- ✓ From atomic- to material-relevant scale
	- ✓ No empiricism
	- ✓ Accurate, transferable
	- ✓ Computable
- ‣ **First-Principles Modeling:** 
	- ‣ Accurate for single component, impractical for material

#### ‣ **Multi-Scale Modeling:**

- ‣ coarse-grained model at large scale
- ‣ most important first-principles, atomistic information ideally retained
- ‣ *balance* of accuracy and computational cost









#### Multi-Scale Modeling of Materials

- ‣ **Perfect** theoretical material model:
	- ✓ From atomic- to material-relevant scale
	- ✓ No empiricism
	- ✓ Accurate, transferable
	- ✓ Computable
- ‣ **First-Principles Modeling:** 
	- ‣ Accurate for single component, impractical for material

#### ‣ **Multi-Scale Modeling:**

- ‣ coarse-grained model at large scale
- ‣ most important first-principles, atomistic information ideally retained
- ‣ *balance* of accuracy and computational cost













#### **Target**: 1st-principles spin dynamics in **realistic molecular semi-conductor** models





#### **Target**: 1st-principles spin dynamics in **realistic molecular semi-conductor** models



ITN SEPOMO | NWE4 Mons, Belgium - Erik R. McNellis 9

# Spin Exchange



- ‣ *Exchange:* neighboring spin inversion, unchanged charge state
- ‣ Molecular wavefunction decays rapidly, exponentially in interstitial region - coupling also
- $\blacktriangleright$  Two-body coupling  $J_{ij} =$  $E_{\uparrow\downarrow}-E_{\uparrow\uparrow}$  $4\langle S_i \rangle \langle S_j \rangle$
- ‣ from constrained DFT1 *E*↑↓, *E*↑↑, ⟨*S*⟩
- ‣ Coupling as function of polaron separation fitted to exponential function2

Spin Exchange





- 1. I. Rudra, Q. Wu and T. Van Voorhis, J. Chem. Phys. 124, 24103 (2006)
- 2. A. R. O'Dea, A. F. Curtis, N. J. B. Green, C. R. Tinunel and P. J. Hore, J. Phys. Chem. A 109, 869 (2005)



![](_page_21_Picture_1.jpeg)

- ‣ *Hyperfine coupling* (HFC): electronic, nuclear spin interaction
	- ‣ vanishes for 'closed-shell' molecules
	- organic elements often nuclear spin free
- ‣ **hydrogens, ionic** molecules main source

(External + Local Hyperfine Field) ⋅ *g*-tensor

![](_page_21_Picture_7.jpeg)

![](_page_21_Picture_8.jpeg)

1. S. Schott, ERM, C. B. Nielsen, H.-Y. Chen, …, J. Sinova, and H. Sirringhaus, Nat. Commun. 8, 15200 (2017). 2. ERM, S. Schott, H. Sirringhaus, and J. Sinova, Phys. Rev. Mater. 2, 074405 (2018)

![](_page_22_Picture_0.jpeg)

# Hyperfine Field (HFI) / *g*-Tensor

- ‣ *Hyperfine coupling* (HFC): electronic, nuclear spin interaction
	- ‣ vanishes for 'closed-shell' molecules
	- organic elements often nuclear spin free
- ‣ **hydrogens, ionic** molecules main source
- ‣ Gyromagnetic coupling ("*g*-") tensor shift: deviation from free electron value
- Depends on spin-orbit coupling (SOC)
- $\blacktriangleright$  Modeling, experiments:<sup>1,2</sup> overlap of electronic spin density with
	- $\triangleright$  nuclear spin  $\rightarrow$  HFI
	- $\rightarrow$  orbital angular momentum  $\rightarrow$ *g*-tensor

![](_page_22_Figure_11.jpeg)

‣ Spin density tuned via **push-pull chemistry** *<sup>g</sup>*-tensor tuning2

![](_page_22_Picture_13.jpeg)

1. S. Schott, ERM, C. B. Nielsen, H.-Y. Chen, …, J. Sinova, and H. Sirringhaus, Nat. Commun. 8, 15200 (2017). 2. ERM, S. Schott, H. Sirringhaus, and J. Sinova, Phys. Rev. Mater. 2, 074405 (2018)

#### (External + Local Hyperfine Field) ⋅ *g*-tensor

![](_page_22_Figure_18.jpeg)

![](_page_23_Picture_1.jpeg)

• Main SOC effect of band transport in traditional semi-conductor materials: scattering between **mixed spin states** – **engleseus** mechanism molecular band Main SOC effect of  $\frac{pana}{hopping}$  transport in spatial Elliott-Yafet<sup>1,2</sup> analogous

- 1.R. J. Elliott, Phys. Rev. 96, 266 (1954)
- 2. Y. Yafet, Adv. Res. Appl. 14, 1 (1963)
- 3. C. A. Masmanidis, H. H. Jaffe, and R. L. Ellis, J. Phys. Chem. 79, 2052 (1975)

![](_page_23_Picture_6.jpeg)

### Spin-Mixing ('Hop-Flip') in Semi-Conductors

![](_page_24_Picture_1.jpeg)

![](_page_24_Figure_2.jpeg)

#### Spin-Mixing ('Hop-Flip') in Semi-Conductors

![](_page_25_Picture_1.jpeg)

![](_page_25_Figure_2.jpeg)

![](_page_26_Picture_1.jpeg)

- $\rightarrow \gamma$  is change in norm of *spin-mixed, perturbed* molecular states  $\langle \psi_0 + | \psi_0 + \rangle = 1 + \gamma^2 = 1 + \gamma_{\uparrow \uparrow}^2 + \gamma_{\uparrow \downarrow}^2$
- ► Spin relaxation  $\alpha$  hopping frequency  $\cdot$   $\gamma^2$

![](_page_26_Picture_4.jpeg)

1. Z. G. Yu, Phys. Rev. B 85, 115201 (2012) 2. U. Chopra, S. Shambhawi, S. A. Egorov, J. Sinova, and ERM, Adv. Func. Mater. (submitted)

- $\rightarrow \gamma$  is change in norm of *spin-mixed, perturbed* molecular states
- ► Spin relaxation  $\alpha$  hopping frequency  $\cdot$   $\gamma^2$
- Original formulation<sup>1</sup>:
- Restricted wavefunction
- $\chi$  Empirical SO constants → minimal basis set ≤ **p**-functions
	- $\rightarrow$  only light molecules, poorly
- ✗ No understanding of xcapproximation influence

1. Z. G. Yu, Phys. Rev. B 85, 115201 (2012) 2. U. Chopra, S. Shambhawi, S. A. Egorov, J. Sinova, and ERM, Adv. Func. Mater. (submitted)

![](_page_27_Picture_10.jpeg)

![](_page_27_Picture_11.jpeg)

 $\langle \psi_0 + | \psi_0 + \rangle = 1 + \gamma^2 = 1 + \gamma_{\uparrow \uparrow}^2 + \gamma_{\uparrow \downarrow}^2$ 

![](_page_27_Picture_13.jpeg)

- $\rightarrow \gamma$  is change in norm of *spin-mixed, perturbed* molecular states  $\langle \psi_0 + | \psi_0 + \rangle = 1 + \gamma^2 = 1 + \gamma_{\uparrow \uparrow}^2 + \gamma_{\uparrow \downarrow}^2$
- ► Spin relaxation  $\alpha$  hopping frequency  $\cdot$   $\gamma^2$
- Original formulation<sup>1</sup>:
- Restricted wavefunction
- $\chi$  Empirical SO constants → minimal basis set ≤ **p**-functions
	- $\rightarrow$  only light molecules, poorly
- ✗ No understanding of xcapproximation influence
- Our generalization:<sup>2</sup>
- ✓ Unrestricted wavefunction
- ✓ Any basis set
- ✓ Any molecule
- ✓ Any single-determinant level of theory, (e.g., DFT)

![](_page_28_Picture_13.jpeg)

1. Z. G. Yu, Phys. Rev. B 85, 115201 (2012) 2. U. Chopra, S. Shambhawi, S. A. Egorov, J. Sinova, and ERM, Adv. Func. Mater. (submitted)

![](_page_28_Picture_17.jpeg)

- $\rightarrow \gamma$  is change in norm of *spin-mixed, perturbed* molecular states  $\langle \psi_0 + | \psi_0 + \rangle = 1 + \gamma^2 = 1 + \gamma_{\uparrow \uparrow}^2 + \gamma_{\uparrow \downarrow}^2$
- ► Spin relaxation  $\alpha$  hopping frequency  $\cdot$   $\gamma^2$
- Original formulation<sup>1</sup>:
- Restricted wavefunction
- $\chi$  Empirical SO constants → minimal basis set ≤ **p**-functions
	- $\rightarrow$  only light molecules, poorly
- ✗ No understanding of xcapproximation influence
- Our generalization:<sup>2</sup>
- ✓ Unrestricted wavefunction
- ✓ Any basis set
- ✓ Any molecule
- ✓ Any single-determinant level of theory, (e.g., DFT)
- $\triangleright$  We have reformulated  $\gamma$  with increased
	- accuracy
	- **•** transferability
- $\triangleright$  ... while maintaining scalability?

1. Z. G. Yu, Phys. Rev. B 85, 115201 (2012)

2. U. Chopra, S. Shambhawi, S. A. Egorov, J. Sinova, and ERM, Adv. Func. Mater. (submitted)

![](_page_29_Picture_21.jpeg)

![](_page_29_Picture_22.jpeg)

#### Effects of Generalization: Model Systems

![](_page_30_Picture_1.jpeg)

- $\triangleright$  New cos<sup>2</sup> dip in  $\gamma^2$  **S**z rotation curve in benzene (cf. (**L**・**S**)2 )
- ‣ More pronounced for stronger SOC (thiophene)

![](_page_30_Figure_4.jpeg)

![](_page_30_Picture_5.jpeg)

#### Effects of Generalization: Model Systems

 $\vec{s}_z$ 

Z

 $L_{\mathsf{Z}}$ 

X

 $\neq$ 

 $^{2}$  [x 10<sup>-8</sup>] (Hole)

 $2.70\frac{1}{0}$ 

0 20 40 60 80 100 120 140 160 180

 $\theta$ 

2.80

2.90

5.30

5.40

7.05

7.20

PW92

PBE0

HF

7.35

Y

![](_page_31_Picture_1.jpeg)

Hole **Electron** ------ 7.35

- $\triangleright$  New cos<sup>2</sup> dip in  $\gamma^2$  **S**z rotation curve in benzene (cf. (**L**・**S**)2 )
- ‣ More pronounced for stronger SOC (thiophene)
- $\triangleright$  Better DFT functionals = large quantitative corrections

![](_page_31_Picture_5.jpeg)

4.89

4.91

4.93

⇥ 2

 $2^{2}$  [x 10<sup>-8</sup>] (Electron)

6.35

6.38

7.20

7.28

### Effects of Generalization: Model Systems

 $S_{Z}$ 

Y

![](_page_32_Picture_1.jpeg)

- $\triangleright$  New cos<sup>2</sup> dip in  $\gamma^2$  **S**<sub>z</sub> rotation curve in benzene (cf. (**L**・**S**)2 )
- ‣ More pronounced for stronger SOC (thiophene)
- $\triangleright$  Better DFT functionals = large quantitative corrections
- $\rightarrow \gamma$  (SOC) depends on relative orientation of π-orbital planes
- $\triangleright$  Effect qualitatively and quantitatively improved for biphenyl twist 1e−07

![](_page_32_Figure_7.jpeg)

![](_page_32_Picture_8.jpeg)

### Accurate Spin Transport in Organic Polymers

![](_page_33_Picture_1.jpeg)

- $\triangleright$  Spin diffusion lengths L<sub>s</sub> of  $\sim$ **1200, 600** (nm) in semicrystalline PBTTT, P3HT polymers1
- ▶ Simple spin diffusion model<sup>2</sup>, generalized  $\gamma$  predicts  $L_s$ within experimental errors

![](_page_33_Figure_4.jpeg)

![](_page_33_Picture_5.jpeg)

1. Wang, Shu-Jen, …, R. Mahani, U. Chopra, ERM, et al, Nat. Electron. 2, 98 (2019) 2. Z. G. Yu, De Gruyter Open: Nanoelectronics and Spintronics 1, 1 (2015)

# Accurate Spin Transport in Organic Polymers

![](_page_34_Picture_1.jpeg)

- $\triangleright$  Spin diffusion lengths L<sub>s</sub> of  $\sim$ **1200, 600** (nm) in semicrystalline PBTTT, P3HT polymers1
- ▶ Simple spin diffusion model<sup>2</sup>, generalized  $\gamma$  predicts  $L_s$ within experimental errors
- Weak other mechanisms, locally high hopping rates  $\rightarrow \gamma$  completely determines L<sub>s</sub>
- $\triangleright$  Variation in  $\gamma$ , L<sub>s</sub> because of **varying π orbital planes** along chain
- ! Want long Ls? **Flatten** your π-conjugated polymer

![](_page_34_Figure_7.jpeg)

![](_page_34_Figure_8.jpeg)

![](_page_34_Picture_9.jpeg)

1. Wang, Shu-Jen, …, R. Mahani, U. Chopra, ERM, et al, Nat. Electron. 2, 98 (2019)

2. Z. G. Yu, De Gruyter Open: Nanoelectronics and Spintronics 1, 1 (2015)

# As a High-Throughput Computational Tool

doi:10.1038/nature12909

![](_page_35_Picture_1.jpeg)

‣ **High-throughput**: characteristic property + robust modeling technique = huge scans of candidate molecules (e.g. batteries<sup>1</sup>, photovoltaics<sup>2</sup>)

#### **LETTER**

A metal-free organic-inorganic aqueous flow battery

Brian Huskinson<sup>1</sup>, Michael P. Marshak<sup>1,2</sup>, Changwon Suh<sup>2</sup>, Süleyman Er<sup>2,3</sup>, Michael R. Gerhardt<sup>1</sup>, Cooper J. Galvin<sup>2</sup>, Xudong Chen<sup>2</sup>, Alán Aspuru-Guzik<sup>2</sup>, Roy G. Gordon<sup>1,2</sup> & Michael J. Aziz<sup>1</sup>

![](_page_35_Picture_7.jpeg)

1. B. Huskinson et al., Nature 505, 195 (2014) 2. J. Hachmann et al., J. Phys. Chem. Lett. 2, 2241 (2011)

![](_page_35_Picture_9.jpeg)

1. S. Schott, U. Chopra, V. Lemaur, A. Melnyk, Yoan Olivier, …, ERM, D. Andrienko, D. Beljonne, J. Sinova, and H. Sirringhaus Nat. Physics. (accepted) 2. U. Chopra, S. A. Egorov, J. Sinova, and ERM, J. Phys. Chem. C, (submitted)

# As a High-Throughput Computational Tool

doi:10.1038/nature12909

![](_page_36_Picture_1.jpeg)

‣ **High-throughput**: characteristic property + robust modeling technique = huge scans of candidate molecules (e.g. batteries<sup>1</sup>, photovoltaics<sup>2</sup>)

#### **LETTER**

A metal-free organic-inorganic aqueous flow battery Brian<mark>[Huskinson<sup>1</sup>\*,</mark>]Michael P. Marshak<sup>1,24</sup>, Changwon Suh<sup>2</sup>, Süleyman Er<sup>2,3</sup>, Michael R. Gerhardt<sup>1</sup>, Cooper J. Galvin<sup>2</sup>,<br>Xudong Chen<sup>2</sup>, Alán Aspuru-Guzik<sup>2</sup>, Roy G. Gordon<sup>1,2</sup> & Michael J. Aziz<sup>1</sup>

- 
- $\triangleright$  Our  $\gamma$  calculation technique
	- ✓ relies on standard DFT
	- $\checkmark$  is highly task parallel
	- $\checkmark$  is highly automatable
- $\rightarrow \gamma$  calculations of **every** state in polymer morphologies possible
- ‣ Statistical picture of polymer spin relaxation otherwise unattainable1,2

![](_page_36_Picture_12.jpeg)

- 1. S. Schott, U. Chopra, V. Lemaur, A. Melnyk, Yoan Olivier, …, ERM, D. Andrienko, D. Beljonne, J. Sinova, and H. Sirringhaus Nat. Physics. (accepted)
	-
- 2. U. Chopra, S. A. Egorov, J. Sinova, and ERM, J. Phys. Chem. C, (submitted)

The Harvard Clean Energy Project: Large-Scale Computational Screening and Design of Organic Photovoltaics on the World **Community Grid** 

PHYSICAL CHEMISTRY Letters

Johannes Hachmann,\*\* Roberto Olivares-Amaya,<sup>†</sup> Sule Atahan-Evrenk,\* Carlos Amador-Bedolla,<sup>†,\*</sup><br>Roel S. Sánchez-Carrera,<sup>|1.†</sup> Aryeh Gold-Parker,<sup>\*</sup> Leslie Vogt,<sup>†</sup> Anna M. Brockway,<sup>5</sup> and Alán Aspuru-Guzik<sup>1,\*</sup>

1. B. Huskinson et al., Nature 505, 195 (2014) 2. J. Hachmann et al., J. Phys. Chem. Lett. 2, 2241 (2011)

![](_page_36_Figure_19.jpeg)

1. K. Bader et al., Chem. Comm. 52, 3623 (2016) 2. U. Chopra, S. Shambhawi, S. A. Egorov, J. Sinova, and ERM, Adv. Func. Mater. (submitted)

ITN SEPOMO | NWE4 Mons, Belgium - Erik R. McNellis 17

 $\gamma$  Transferable to Metal-Phthalocyanines (MPcs)

- $\triangleright$  Bader<sup>1</sup> et al. measure longitudinal spin relaxation times  $T_1$  in dissolved MPcs
- ‣ MPcs too complex for old formulation

![](_page_37_Figure_6.jpeg)

![](_page_37_Picture_7.jpeg)

![](_page_37_Picture_8.jpeg)

ITN SEPOMO | NWE4 Mons, Belgium - Erik R. McNellis 17

#### Transferable to Metal-Phthalocyanines (MPcs)

- $\triangleright$  Bader<sup>1</sup> et al. measure longitudinal spin relaxation times  $T_1$  in dissolved MPcs
- ‣ MPcs too complex for old formulation
- $T_1 \propto$ If spin relaxes through charge hopping,
- ► Fit of  $T_1 \approx \kappa / \gamma^2$  predicts experiment to
	- **~ 40 %** over **4 orders** of magnitude

1. K. Bader et al., Chem. Comm. 52, 3623 (2016) 2. U. Chopra, S. Shambhawi, S. A. Egorov, J. Sinova, and ERM, Adv. Func. Mater. (submitted)

![](_page_38_Figure_9.jpeg)

![](_page_38_Picture_10.jpeg)

 $M = VO$ , Mn, Co, Cu

ITN SEPOMO | NWE4 Mons, Belgium - Erik R. McNellis 17

▶ Bader<sup>1</sup> et al. measure longitudinal spin relaxation times  $T_1$  in dissolved MPcs

‣ MPcs too complex for old formulation

- If spin relaxes through charge hopping,
- ► Fit of  $T_1 \approx \kappa / \gamma^2$  predicts experiment to **~ 40 %** over **4 orders** of magnitude
- ! Spin relaxation indistinguishable mix of (maybe) hopping and **thermal** effects
- also  $\propto \hat {\mathrm{H}}_{\mathrm{SOC}}^2$  fit still works  $\sqrt{\ }$  Thermal (spin-phonon coupling)
- ✓ SOC highly accurate

 $\triangleright$  Can we do equally well with thermal effects with method to  $\gamma^2$  standard?

1. K. Bader et al., Chem. Comm. 52, 3623 (2016)

2. U. Chopra, S. Shambhawi, S. A. Egorov, J. Sinova, and ERM, Adv. Func. Mater. (submitted)

![](_page_39_Figure_11.jpeg)

![](_page_39_Figure_12.jpeg)

 $=$  VO, Mn, Co, Cu

![](_page_39_Picture_13.jpeg)

# Spin-Phonon Coupling

![](_page_40_Picture_1.jpeg)

- ‣ SOC coupling to a virtual state of opposite spin, resonant with zero-field Zeeman split via electron-phonon coupling
- ‣ Generally requires multiple phonons for resonance
- ‣ *Elastic*: spin relaxes on isolated, thermally excited molecule
- ‣ *Inelastic*: phonon absorption / emission

quantum dots<sup>1</sup>, defects in solids<sup>2</sup>

![](_page_40_Figure_6.jpeg)

Absorption followed by emission (or vice versa) via virtual state

- 
- 1. Y. G. Semenov and K. W. Kim, Phys. Rev. B 75, 195342 (2007) 2. S. A. Egorov and J. L. Skinner, J. Chem. Phys. 103, 1533 (1995) 3. S. Roychoudhari, S. Sanvito, PRB, 98, 125204 (2018)

 $\triangleright$  Updated: adapted from crystals<sup>3</sup>, same SOC as  $\gamma$ 

![](_page_41_Picture_1.jpeg)

**Target**: 1st-principles spin dynamics in **realistic molecular semi-conductor** models

![](_page_41_Figure_3.jpeg)

ITN SEPOMO | NWE4 Mons, Belgium - Erik R. McNellis

![](_page_42_Picture_1.jpeg)

**Target**: 1st-principles spin dynamics in **realistic molecular semi-conductor** models

![](_page_42_Figure_3.jpeg)

ITN SEPOMO | NWE4 Mons, Belgium - Erik R. McNellis

# VOTCA-STP

![](_page_43_Picture_1.jpeg)

‣ Our approach: spin-dynamics **on top** of multi-scale charge-dynamics

VOTCA-CTP1:

- ‣ Hopping charge transport in soft matter
- $\triangleright$  MD, Marcus theory  $\rightarrow$  thermal effects

![](_page_43_Figure_6.jpeg)

![](_page_43_Figure_7.jpeg)

![](_page_43_Picture_8.jpeg)

1. V. Rühle, A. Lukyanov, F. May et al, J. Chem. Theory Comput. 7, 3335 (2011) 2. Z. G. Yu, F. Ding and H. Wang, Phys. Rev. B 87, 205446 (2013) 3. U. Chopra, S. Shambhawi,…, and ERM, Adv. Func. Mater. (submitted) 4. I. Rudra, Q. Wu, T. Van Voorhis, Inorg. Chem. 46, 10539 (2007)

![](_page_43_Picture_10.jpeg)

# VOTCA-STP

![](_page_44_Picture_1.jpeg)

‣ Our approach: spin-dynamics **on top** of multi-scale charge-dynamics

VOTCA-CTP1:

- ‣ Hopping charge transport in soft matter
- $\triangleright$  MD, Marcus theory  $\rightarrow$  thermal effects

'VOTCA-**STP**':

- ► Separate spin dynamics KMC loop
- ‣ Single-site mechanisms:
	- ‣ hyperfine field2 / *g*-tensor
	- ‣ thermal relaxation
- ‣ Two-site mechanisms:
	- $\triangleright$  spin-flip at hop (rate from  $\gamma^3$ )
	- $\blacktriangleright$  spin exchange<sup>4</sup>
	- ▶ spin dipole

#### VOTCA-CTP + Spin Dynamics

![](_page_44_Figure_16.jpeg)

![](_page_44_Picture_17.jpeg)

1. V. Rühle, A. Lukyanov, F. May et al, J. Chem. Theory Comput. 7, 3335 (2011)

- 2. Z. G. Yu, F. Ding and H. Wang, Phys. Rev. B 87, 205446 (2013)
- 3. U. Chopra, S. Shambhawi,…, and ERM, Adv. Func. Mater. (submitted)
- 4. I. Rudra, Q. Wu, T. Van Voorhis, Inorg. Chem. 46, 10539 (2007)

![](_page_44_Picture_22.jpeg)

![](_page_44_Picture_23.jpeg)

# VOTCA-STP Spin Dynamics Concept

![](_page_45_Picture_1.jpeg)

- ‣ Material model randomly seeded with charges
- ▶ Changes in spin polarization monitored as charges move
- ‣ **Explicit link** between chargeand spin-dynamics
- ‣ Allows for **unprecedented inference** of one from the other

![](_page_45_Figure_6.jpeg)

Spin Polarization Change, Current

![](_page_45_Picture_8.jpeg)

# *T*1 Spin Relaxation Time in Bulk Alq3

- $\triangleright$  Proof of concept: bulk Alq<sub>3</sub> longitudinal spin relaxation time
- ‣ Amorphous bulk model: 4096 molecule cell, periodic boundary conditions, monopolaronic transport approximation

![](_page_46_Picture_3.jpeg)

![](_page_46_Picture_6.jpeg)

![](_page_46_Figure_7.jpeg)

# *T*1 Spin Relaxation Time in Bulk Alq3

- $\triangleright$  Proof of concept: bulk Alq<sub>3</sub> longitudinal spin relaxation time
- ‣ Amorphous bulk model: 4096 molecule cell, periodic boundary conditions, monopolaronic transport approximation

![](_page_47_Picture_3.jpeg)

![](_page_47_Picture_6.jpeg)

![](_page_47_Figure_7.jpeg)

### *T*1 Spin Relaxation Time in Bulk Alq3

- $\triangleright$  Proof of concept: bulk Alq<sub>3</sub> longitudinal spin relaxation time
- ‣ Amorphous bulk model: 4096 molecule cell, periodic boundary conditions, monopolaronic transport approximation
- ‣ Single molecular dynamics snapshot from glassy phase of bulk Alq3
- ‣ Large internal variations in molecular …
	- internal geometry
	- $\triangleright$  density
	- $\blacktriangleright$  relative orientation
- ... lead to large variations in
	- ‣ single molecule SOC
	- ‣ local hyperfine fields
	- ‣ current density
- 1. V. Rühle, A. Lukyanov, F. May et al, J. Chem. Theory Comput. 7, 3335 (2011)

current density isocontour1

![](_page_48_Picture_15.jpeg)

![](_page_48_Figure_16.jpeg)

![](_page_48_Picture_17.jpeg)

ITN SEPOMO | NWE4 Mons, Belgium - Erik R. McNellis 22

1e21

#### **• thermal relaxation negligible**

‣ With **no exchange**, 'hockey-stick'  $T_1$  shape:

- ► HFI and hop-flipping interfere, slowing  $T_1$
- ‣ Charge hopping blocked as concentration increases, slower HFI relaxation dominates

![](_page_49_Picture_6.jpeg)

![](_page_49_Picture_7.jpeg)

![](_page_49_Picture_8.jpeg)

 $\triangleright$  Proof of concept: bulk Alq<sub>3</sub> longitudinal spin relaxation time

boundary conditions, monopolaronic transport approximation

300

‣ Amorphous bulk model: 4096 molecule cell, periodic

![](_page_49_Picture_10.jpeg)

*T*1 Spin Relaxation Time in Bulk Alq3

![](_page_49_Picture_11.jpeg)

ITN SEPOMO | NWE4 Mons, Belgium - Erik R. McNellis 22

#### ‣ With **exchange included,** large speedup in high concentration regime

- ‣ Due to relaxation **traps**:
- ‣ Exchange-mediated spin diffusion enhances access to fast relaxation sites
- ‣ Complex balance of effects

![](_page_50_Figure_5.jpeg)

- $\triangleright$  Proof of concept: bulk Alq<sub>3</sub> longitudinal spin relaxation time
- ‣ Amorphous bulk model: 4096 molecule cell, periodic boundary conditions, monopolaronic transport approximation

*T*1 Spin Relaxation Time in Bulk Alq3

![](_page_50_Picture_10.jpeg)

![](_page_50_Picture_11.jpeg)

#### **Summary**

![](_page_51_Picture_1.jpeg)

- ‣ Fully first-principles modeling of molecular and organic semi-conductor materials **possible**
- ‣ Theoretical modeling offers **otherwise unattainable** insights through
	- ‣ **versatility** at all scales
	- ‣ **complementarity** to experiment
- ‣ Models highlight **complexity** of molecular spin dynamics - exciting!
- ‣ Still plenty of work to do to raise
	- **accuracy**
	- **•** transferability
	- ‣ scalability
- ‣ … of current methods

![](_page_51_Picture_12.jpeg)

#### **Summary**

![](_page_52_Picture_1.jpeg)

- ‣ Fully first-principles modeling of molecular and organic semi-conductor materials **possible**
- ‣ Theoretical modeling offers **otherwise unattainable** insights through
	- ‣ **versatility** at all scales
	- ‣ **complementarity** to experiment
- ‣ Models highlight **complexity** of molecular spin dynamics - exciting!
- ‣ Still plenty of work to do to raise
	- **accuracy**
	- **•** transferability
	- ‣ scalability
- ‣ … of current methods

#### **Thank You!**

![](_page_52_Picture_13.jpeg)